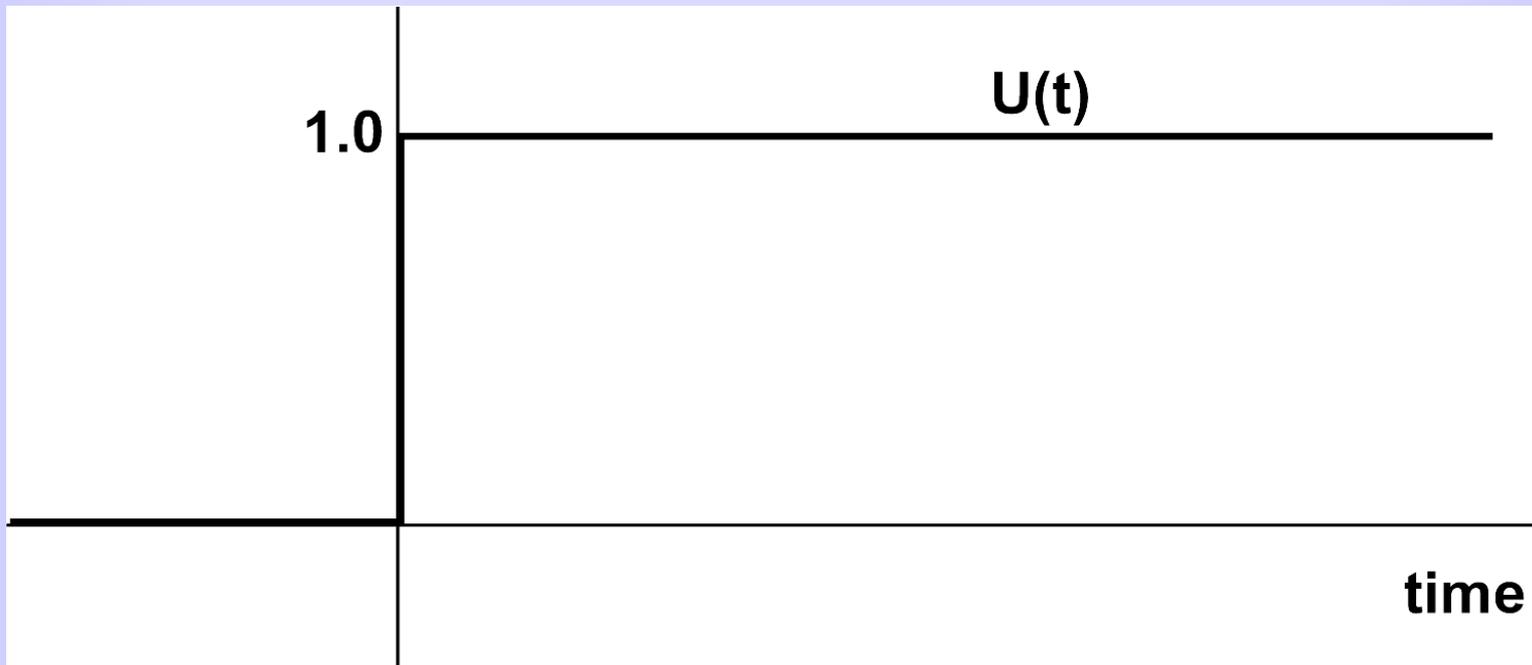


Laplace Transformation, Pole-Zero analysis, the Oculomotor Plant and the VOR, revisited

Problem: the Fourier transform of many, often used functions is not defined! E.g., the **step function**:



$$F(\omega) = \int_{-\infty}^{\infty} U(t)e^{-j\omega t} dt = \int_0^{\infty} e^{-j\omega t} dt = -\frac{1}{j\omega} e^{-j\omega t} \Big|_0^{\infty} = ??$$

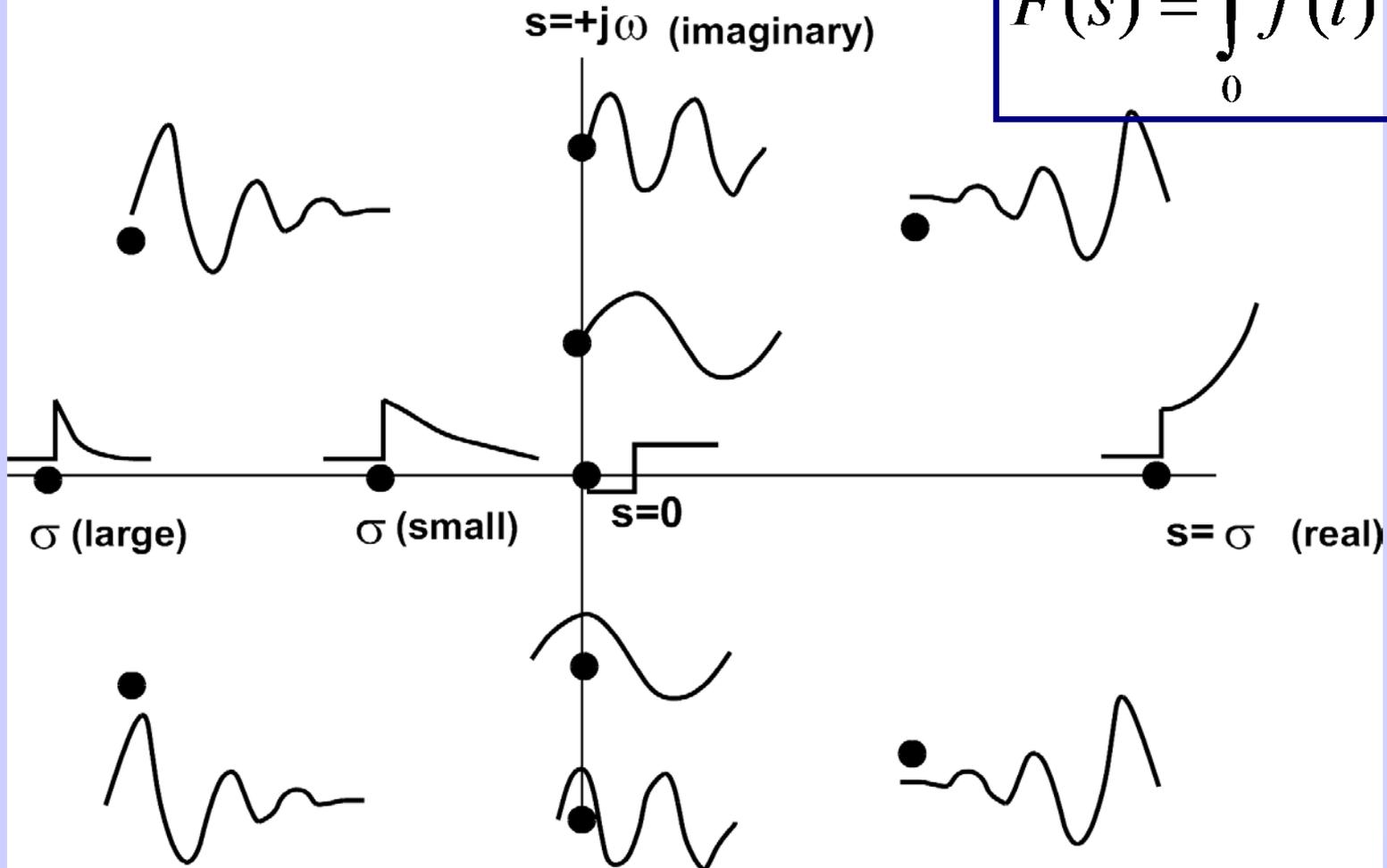
'weighting function'

Laplace transform: uses a new weighting function!

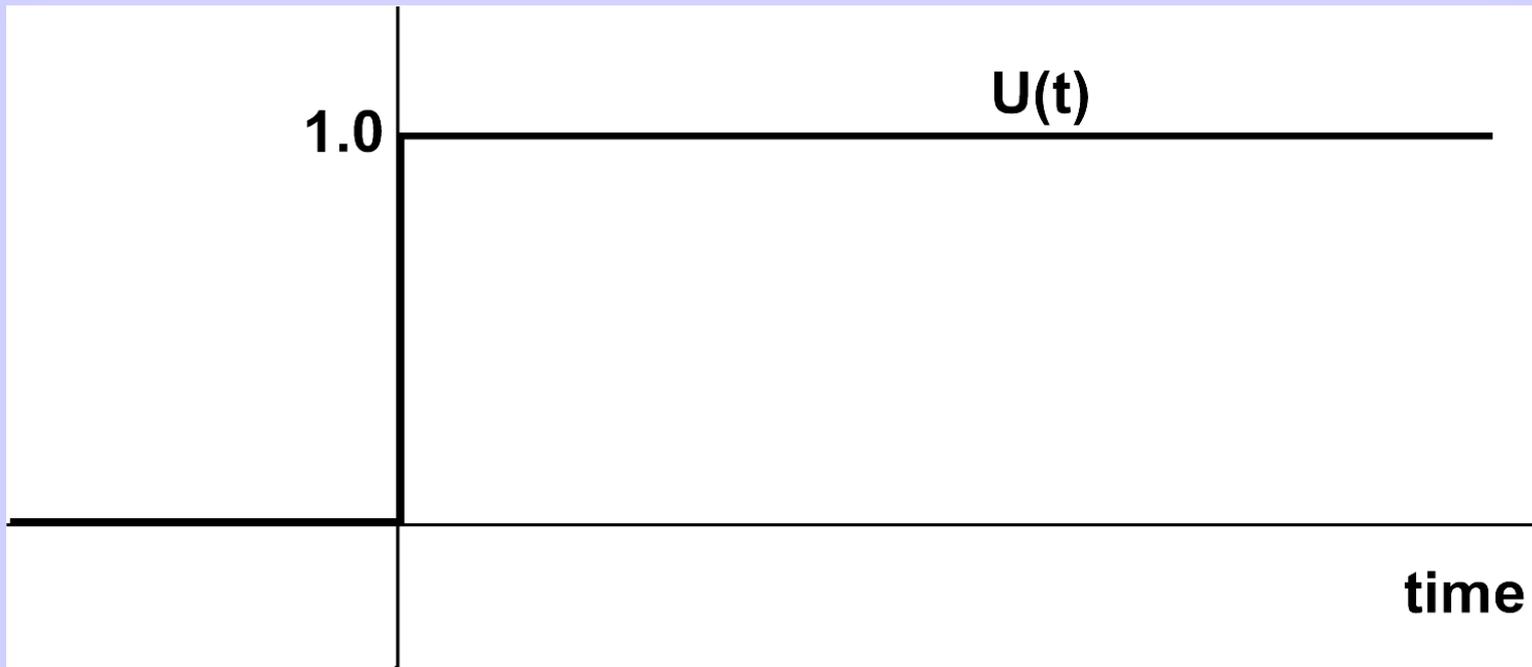
$$e^{-st} \quad s \equiv \sigma + j\omega$$

The Complex s-plane

$$F(s) = \int_0^{\infty} f(t) \cdot e^{-st} dt$$



The Laplace transform for the **step function**:



$$F(s) = \int_0^{\infty} U(t) \cdot e^{-st} dt = \int_0^{\infty} e^{-st} dt = \frac{1}{s}$$

If $f(t)$ en $g(t)$ \rightarrow $F(s)$ en $G(s)$

then (left as Exercises for you):

$a \cdot f(t)$	\rightarrow	$a \cdot F(s)$	scaling
 $f(t \pm a)$	\rightarrow	$\exp(\pm a \cdot s) \cdot F(s)$	delay
 df/dt	\rightarrow	$s \cdot F(s)$	derivative
 $\int f(t) \cdot dt$	\rightarrow	$F(s)/s$	integral
 $\exp(\pm a \cdot t)$	\rightarrow	$1/(s \mp a)$	exponential
$U(t)$	\rightarrow	$1/s$	step
$\delta(t)$	\rightarrow	1	Dirac pulse
 $\exp(\pm a \cdot t) \cdot f(t)$	\rightarrow	$F(s \mp a)$	exp. decay
$h(t) = \int f(x) \cdot g(t-x) \cdot dx$	\rightarrow	$H(s) = F(s) \cdot G(s)$	convolution

We apply the Laplace transform to the (1st-order) model of the eye muscles (the 'oculomotor plant'):

$$\Delta R_m(t) = k \cdot E(t) + r \cdot \frac{dE}{dt}$$

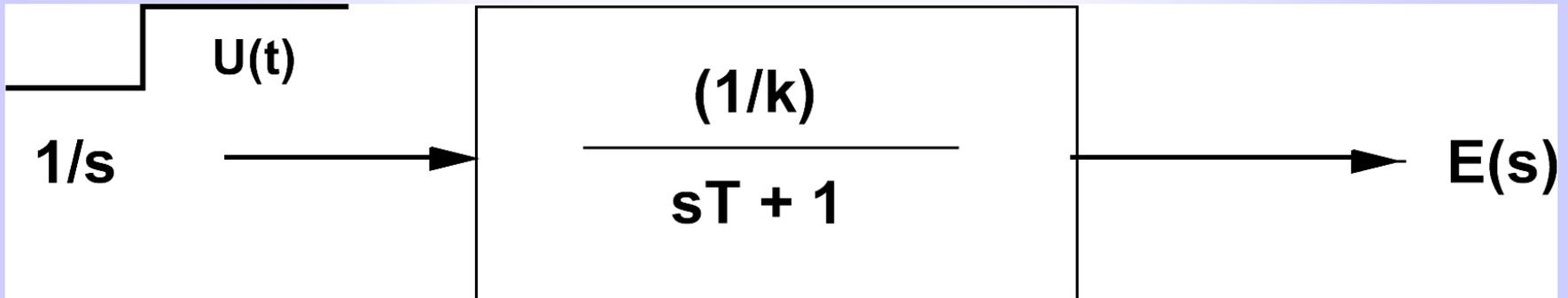
Laplace Transformation:

$$\Delta R_m(s) = k \cdot E(s) + r \cdot s \cdot E(s) = (k + r \cdot s) \cdot E(s)$$

The Laplace transfer function (output/input) of the plant is then:

$$H(s) \equiv \frac{E(s)}{\Delta R_m(s)} = \frac{1}{k + r \cdot s} = \frac{(1/k)}{1 + s \cdot (r/k)} = \frac{(1/k)}{1 + s \cdot T_E}$$

Example:
 Compute the **Step response** for the 1st-order plant:



$$E(s) = \frac{1}{s} \cdot \frac{(1/k)}{1 + sT} = \frac{(1/k)}{s} - \frac{(1/k)}{s + (1/T)} \xrightarrow{\text{LT}^{-1}} E(t) = \frac{1}{k} \cdot (1 - e^{-t/T})$$

Important: first rewrite the Laplace function into its standard form before applying LT^{-1} !

Interestingly, we can get our results even faster!

For this, we only need to look at the **poles** and **zeros** of the LT:

$$H(s) \equiv \frac{\text{numerator}}{\text{denominator}} = \frac{N(s)}{D(s)} \quad \begin{array}{l} \text{Zeros: } N(s) = 0 \\ \text{Poles: } D(s) = 0 \end{array}$$

Because $N(s)$ en $D(s)$ are complex functions, and $H(s)$ is derived from an arbitrary order linear differential equation, it follows that $N(s)$ and $D(s)$ can **always** be written as the product of zero-, first-, and second-order polynomials in s :

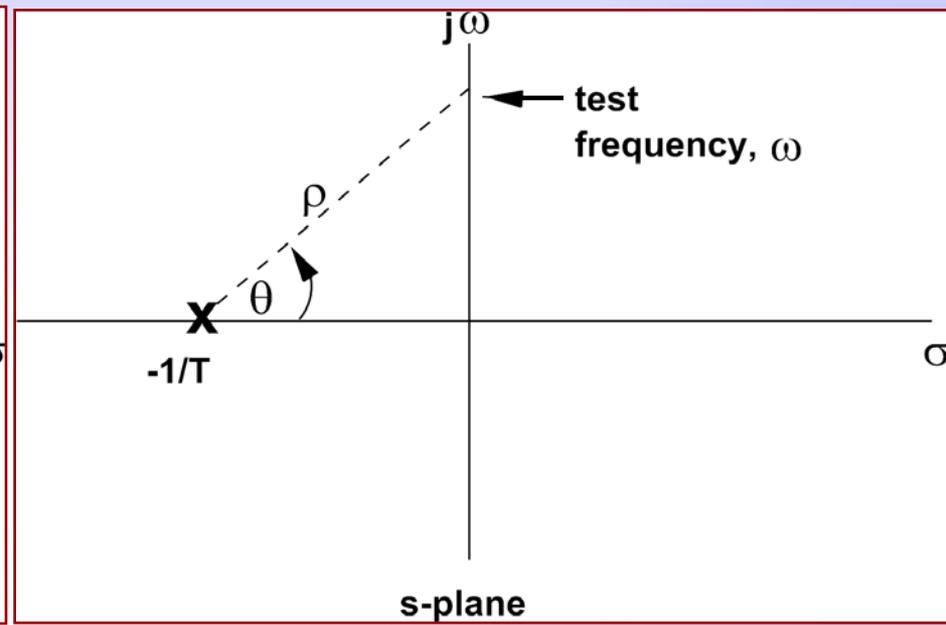
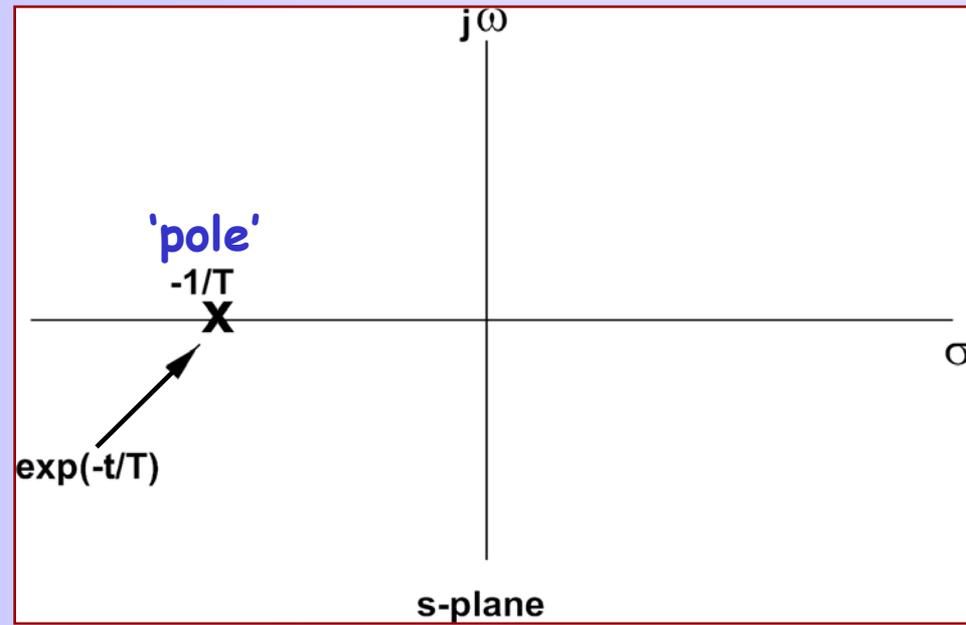
$$N(s) = \sum_{n=1}^N a_n \cdot s^n = (a + b \cdot s) \cdot (c + d \cdot s + e \cdot s^2) \cdot \dots$$

The 1st-order polynomials yield zeros or poles that lie on real (σ -)axis of s -space (i.e., exponential functions of time)
The 2nd-order polynomials **always** yield complex-conjugated pairs as factors: $s_{1,2} = a \pm j \cdot \omega$

How does this work for our simple eye-plant model?

Impulse response of the plant:

$$H(s) \equiv \frac{E(s)}{\Delta R_m(s)} = \frac{(1/k)}{1 + s \cdot T_E}$$



Reading Gain- and Phase characteristics:

From the location of the pole, we can immediately write the impulse response function in the time domain, as well as the gain- and phase characteristics!

$$G(\omega) = \frac{1}{\rho}$$

$$\Phi(\omega) = -\theta$$

In general the following holds for more complex transfer functions:

- 1) Amplitude characteristics multiply with each other
- 2) Phase characteristics are added to each other (lag: -, advance: +).
- 3) Amplitude characteristic for a pole: $1/\rho(\omega)$, and for a zero: $\rho(\omega)$
- 4) Phase characteristic for a pole: negative angle $\theta(\omega)$ (phase lag)
for a zero: positive angle $\theta(\omega)$ (phase advance)

For example:

than the zero is: $s_1 = -1/a$

And two poles: $s_2 = -1/b$ and $s_3 = -1/c$

$$H(s) \equiv \frac{1 + a \cdot s}{(1 + b \cdot s) \cdot (1 + c \cdot s)}$$

In the frequency domain the transfer characteristic then becomes:

$$G(\omega) = \frac{\sqrt{\omega^2 + (1/a)^2}}{\sqrt{\omega^2 + (1/b)^2} \cdot \sqrt{\omega^2 + (1/c)^2}} \quad \text{and}$$

$$\Phi(\omega) = \arctan(a \cdot \omega) - \arctan(b \cdot \omega) - \arctan(c \cdot \omega)$$

Time behaviour of $h(t)$ is found by inverse LT (after fraction splitting....): at WC!

Model analysis through Laplace transformation:

Some examples:

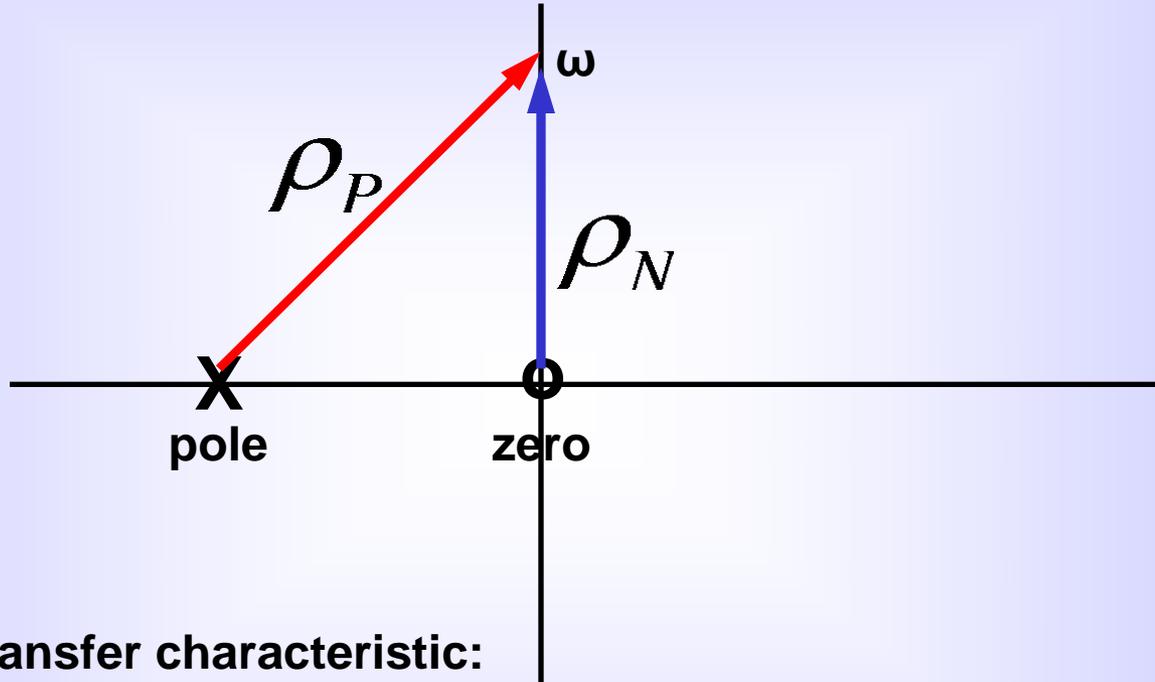
- the VOR
- 'velocity storage'
- neural integration

the VOR:

$$\frac{d\dot{H}}{dt} = \frac{1}{T_{vor}} \cdot \dot{E}(t) + \frac{d\dot{E}}{dt}$$

LT

$$T(s) \equiv \frac{\dot{E}}{\dot{H}} = \frac{sT_{vor}}{1 + sT_{vor}}$$

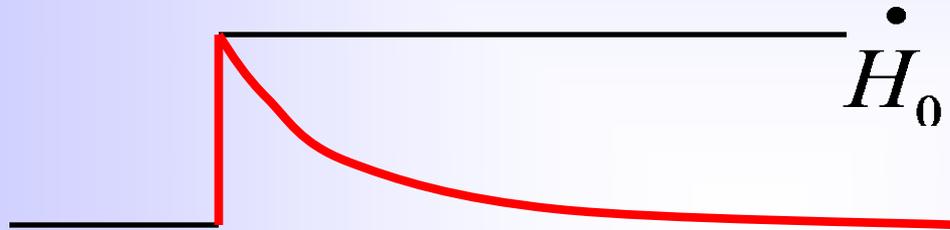


$$G_{vor}(\omega) = G_P(\omega) \cdot G_Z(\omega) = \frac{\rho_Z(\omega)}{\rho_P(\omega)} = \frac{\omega}{\sqrt{\omega^2 + (1/T)^2}}$$

$$\Phi(\omega) = \Phi_Z(\omega) - \Phi_P(\omega) = (\pi/2) - \arctan(\omega T)$$

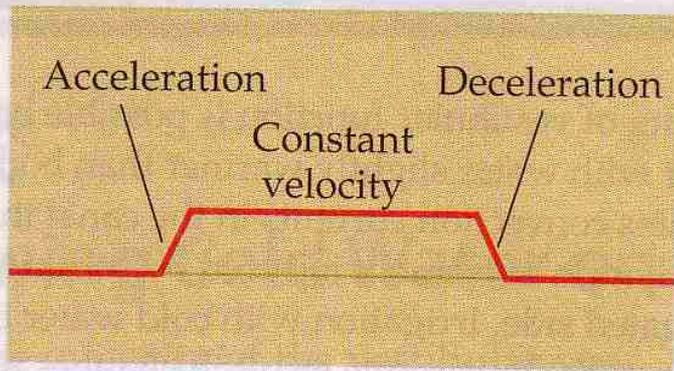
Step response of the VOR?

$$\dot{E}(s) = \frac{\dot{H}_0}{s} \times \frac{sT_{VOR}}{1 + sT_{VOR}} = \frac{\dot{H}_0 T_{VOR}}{1 + sT_{VOR}}$$

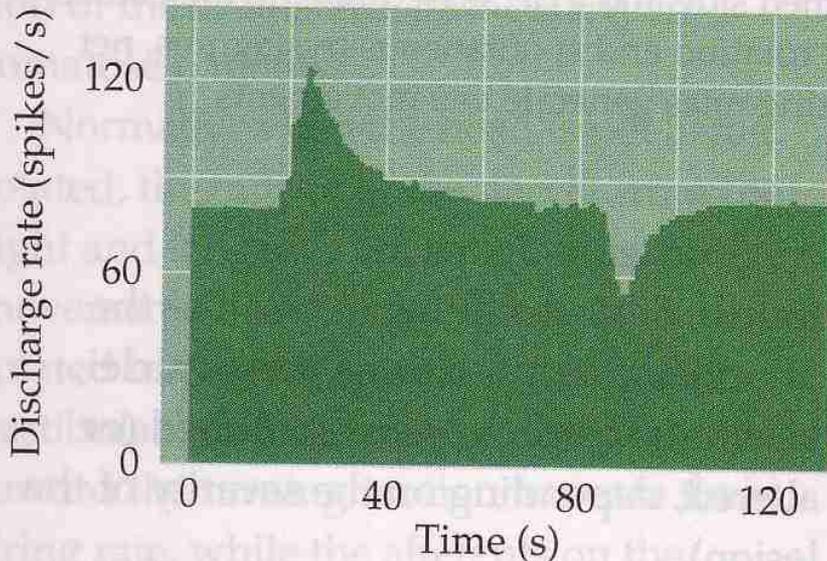


$$\dot{E}(t) = \dot{H}_0 \cdot e^{-t/T_{VOR}}$$

with $T_{VOR} \approx 21$ s



There appears to be an interesting problem in relation to the VOR:



$$R_{vest}(t) = \dot{H}_0 \cdot e^{-t/T_{can}}$$

with $T_{can} \approx 7 \text{ s}$

Recording of neural activity from the vestibular nerve (from Purves, Ch 13) during a constant-velocity head rotation